

Self Induction: It is the property of the coil which opposes any changes in current passing through it.

co-efficient of self induction: L (self induction)

$$\phi \propto I$$

$$\phi = L I \quad \text{where } L \rightarrow \text{coeff. of self induction.}$$

If $I = 1A$ then $\boxed{\phi = L}$

The self inductance of a coil is numerically equal to the mag. flux linked with it when a current of 1A flows through it.

$$\frac{d\phi}{dt} = L \frac{dI}{dt}$$

$$\therefore e = -L \frac{dI}{dt} \quad \text{If } \frac{dI}{dt} = 1 \text{ A s}^{-1} \text{ then}$$

$$\boxed{e = L}$$

\therefore self inductance of a coil is numerically equal to emf induced in the coil when current in it changes at the rate of 1 A s^{-1}

\therefore unit of self inductance is henry (H)

Define henry: $e = L \frac{dI}{dt}$

$$\text{If } \frac{dI}{dt} = 1 \text{ A s}^{-1}, e = 1 \text{ volt then}$$

$$\boxed{L = 1 \text{ H}}$$

The self inductance of a coil is said to be 1 henry when a current in it changing at the rate of 1 A s^{-1} , induces an emf of 1V in the coil.

Expression for Self inductance of a coil:

Consider a coil of length l , area A and turns N . Let I be the current flowing then mag. field is given by $B = \frac{\mu_0 NI}{l}$ tesla — (1)

Total flux linked with the coil $\phi = NBA$

$$\phi = N \left(\frac{\mu_0 NI}{l} \right) A = \frac{\mu_0 N^2 I A}{l} \quad \text{--- (2)}$$

By definition $\phi = LI$ — (3)

where $L \rightarrow$ self inductance of the coil

Equating eqns (2) & (3)

$$L = \frac{\mu_0 N^2 A}{l} \text{ henry}$$

where $\mu_0 = 4\pi \times 10^{-7} \text{ Hm}^{-1}$

If the coil is wound on material of relative permeability μ_r then

$$L = \frac{\mu_0 \mu_r N^2 A}{l}$$

Energy stored in a coil:

Consider a coil of inductance ' L ' in a circuit, let the current grow from zero to maximum value I_0

The induced emf $\bar{E} = -L \frac{dI}{dt}$

This induced emf tends to prevent the growth of current. To maintain the growth,

work has to be done against the back emf.
The work done in small interval of time dt is

$$dW = -E I dt$$

$$dW = -\left[-L \frac{dI}{dt}\right] I dt$$

$$dW = L I dI$$

Total work done $W = \int_0^{I_0} L I dI$

$$W = \left[\frac{L I^2}{2}\right]_0^{I_0}$$

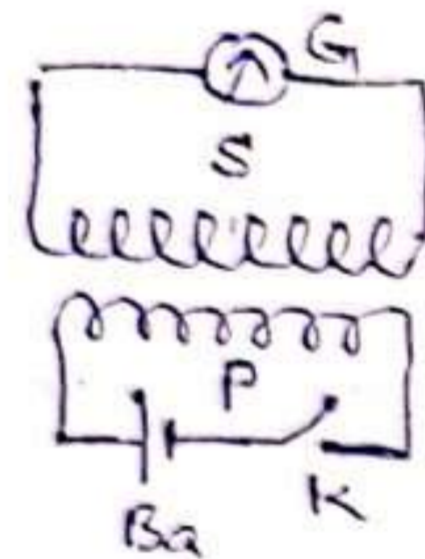
$$W = \frac{1}{2} L I_0^2$$

This work done is stored as energy in the form of magnetic field \therefore

$$E = \frac{1}{2} L I_0^2$$

Mutual induction:-

Consider a secondary coil S placed in the mag. field of the primary coil P. when the current



in the coil P increases, mag flux linked with coil S also increases and vice versa. The galvanometer shows deflection. This shows that an emf and current is induced in coil S due to change of current in the coil P.

\therefore Mutual induction is the phenomenon of inducing emf in one coil due to current changes in a neighbouring coil.

Mutual inductance or coefficient of mutual induction:-

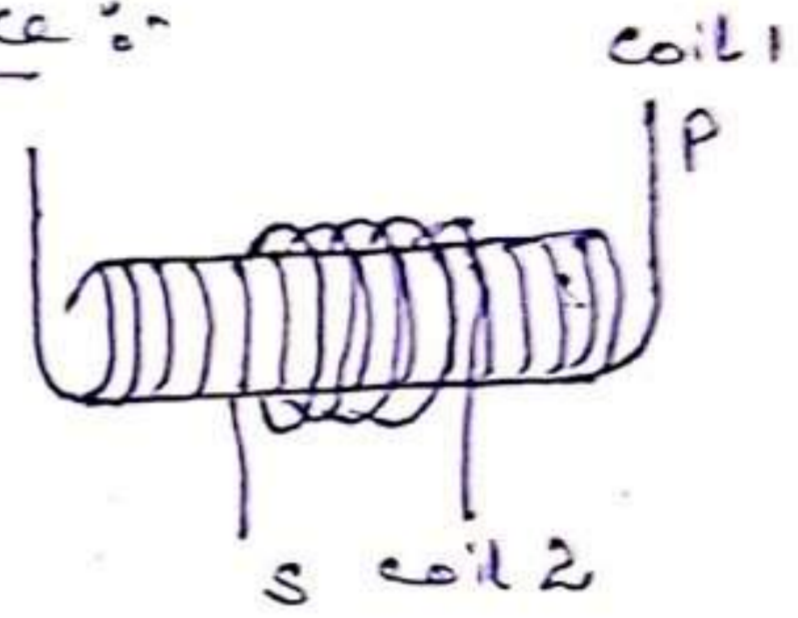
Consider $e = M \frac{di}{dt}$

when $\frac{di}{dt} = 1 \text{ A s}^{-1}$ then $M = e$

Mutual inductance between pair of coils is numerically equal to the emf induced in one coil when current changes at the rate of one ampere per second in the neighbouring coil. S.I. unit of mutual inductance is henry (H)

Expression for mutual inductance:-

Consider a primary coil P of length L, area A and no of turns N_p , consider a short secondary coil S of turns N_s wound closely over central portion of coil P.



If I be the current flowing in P then mag. field

$$B = \frac{\mu_0 N_p I}{L}$$

Flux linked in Primary P is $\phi' = BA$

$$\phi' = \frac{\mu_0 N_p I A}{L} \quad \text{--- (1)}$$

Total flux linked in Secondary S is $\phi = N_s \phi'$ --- (2)

$$\phi = \frac{\mu_0 N_p N_s I A}{L} \quad \text{--- (3)}$$

But $\phi = MI$ --- (4) comparing eqy (3) & (4)

$$M = \frac{\mu_0 N_p N_s A}{L}$$

If the coil is wound material of relative permeability

then

$$M = \frac{\mu_0 \mu_r N_p N_s A}{L}$$

Eddy current:- when a metal block like copper disc is in relative motion with a magnetic field, induced currents are set up in it, opposing the relative motion. These induced currents are called eddy currents.

Applications:

- 1) Electromagnetic damping: In a moving coil galvanometer the coil is wound on a metal frame. when the frame rotates with the coil in the mag. field, eddy currents are formed in it opposing its motion and thus the coil is quickly brought to rest.
- 2) Induction furnace: A high frequency alternating current is passed through a circuit which surrounds the metal furnace the eddy currents are produced in the metal produce high temp.
- 3) Electric brakes: The axle of train is fitted with metal drums. These drums rotate along with the axle. when train is to be stopped, a strong magnetic field is applied on these drums using electromagnet. The eddy current generated in the drums opposes the motion of axle. Hence the train is brought to rest.

4) Speedometer :- A magnet is placed inside an aluminium drum and is carefully pivoted and held in position by hair spring. This magnet revolves according to the speed of the vehicle. Eddy currents in the drum are set up and drag it through an angle which depends on the speed of the vehicle.

5) Induction motor : A metallic cylinder is placed in a rotating magnetic field and the eddy current produced in the metallic cylinder decreases the relative motion between the cylinder and the rotating mag. field. As the mag. field is continuously rotated metallic cylinder is set up in rotation due eddy current.

Disadvantages of Eddy currents :

- 1) Eddy currents dissipate electrical energy in the form of heat in transformer, dynamo etc.
- 2) Heat produced due to eddy currents breaks the insulation in the electric circuits.

Derivation for paramagnetic susceptibility (1) Curie's law

(2) Langevin's theory of paramagnetism:

Paramagnetic substances are characterized by a molecules having a permanent magnetic moments. As the molecules have permanent dipole moment, they do get aligned in an external field.

Consider a paramagnetic substance having 'N' molecules atoms per unit volume at a temperature 'T'. Each atom/molecule is assumed to possess a permanent magnetic moment \vec{m} . If all the dipoles are along the applied magnetic field, then the total magnetization (\vec{I}) of all dipoles will be $N\vec{m}$. The magnetic potential energy of the dipole which is inclined at an angle θ with the applied field \vec{B} is given by

$$U = -\vec{m} \cdot \vec{B} = -mB \cos \theta \quad \rightarrow (1)$$

From Boltzmann distribution law, the relative probability of finding a dipole making an angle θ with the applied \vec{B} is given by

$$e^{-U/KT} \quad \text{where } K = \text{Boltzmann Constant.}$$

Let dN be the no. of dipoles inclined b/w the angles θ and $\theta + d\theta$ with the field \vec{B} .

$$\text{Then } dN \propto e^{-U/KT} d\Omega$$

$$dN = A e^{-U/KT} d\Omega \quad \rightarrow (2)$$

where 'A' is a constant of proportionality and $d\Omega$ is the solid angle subtended by the cones of semi-vertical angle

θ and $\theta + d\theta$ and the solid angle is given by $d\Omega = 2\pi \sin \theta d\theta$

$$\therefore dN = A e^{-U/KT} \cdot 2\pi \sin \theta d\theta$$

$$\text{Using Equation (1) } dN = A_0 (mB \cos \theta) / KT \cdot 2\pi \sin \theta d\theta$$

$$dN = 2\pi A_0 e^{-\alpha \cos \theta} \sin \theta d\theta \quad \rightarrow (3)$$

$$\text{here } \alpha = mB/KT$$

Total no. of dipoles per unit volume (N) can be obtained by

(1)

Integrating the Equation (1) Now the limits '0' and ' π '
 i.e. $N = \int_0^\pi dN = \int_0^\pi d\theta A e^{\alpha \cos \theta} \sin \theta d\theta$

$$N = 4\pi A \int_0^\pi e^{\alpha \cos \theta} \sin \theta d\theta$$

$$\text{But } \int_0^\pi e^{\alpha \cos \theta} \sin \theta d\theta = \frac{\sinh \alpha}{\alpha}$$

$$\therefore N = \frac{4\pi A \sinh \alpha}{\alpha}$$

(2)

$$A = \frac{N\alpha}{4\pi \sinh \alpha} \rightarrow (4)$$

Using the Value of 'A' in Equation (3) we get

$$dN = \frac{N\alpha}{4\pi \sinh \alpha} (e^{\alpha \cos \theta}) \sin \theta d\theta \rightarrow (5)$$

The Component of Each dipole moment parallel to applied field \vec{H} is $m \cos \theta$, hence the total magnetic moment per unit volume (or total magnetization) of the substance is given by

$$I = \int_0^\pi m \cos \theta dN$$

Using Equation (5)

$$I = \frac{mN\alpha}{4\pi \sinh \alpha} \int_0^\pi e^{\alpha \cos \theta} \sin \theta \cos \theta d\theta$$

$$I = \frac{mN\alpha}{4\pi \sinh \alpha} \left[\frac{e^{\alpha \cos \theta}}{\alpha} - \frac{\cos \theta e^{\alpha \cos \theta}}{\alpha} \right]_0^\pi$$

$$I = \frac{mN}{4\pi \sinh \alpha} \left[\frac{e^{\alpha \cos \theta}}{\alpha} - \cos \theta e^{\alpha \cos \theta} \right]_0^\pi$$

$$I = \frac{mN}{4\pi \sinh \alpha} \left[\frac{e^{-\alpha}}{\alpha} + e^{-\alpha} - \frac{e^{\alpha}}{\alpha} + e^{\alpha} \right]$$

on simplifying above Equation, we get

$$I = Nm \left(\coth \alpha - \frac{1}{\alpha} \right) \rightarrow (6)$$

Since Nm is the Saturation value of the magnetic moment.

$$\text{hence } Nm = I_0$$

i.e., when the substance is completely aligned with the applied field then the magnetic moment is known as saturation magnetic moment (I_0)

$$\therefore I = I_0 \left(\coth \alpha - \frac{1}{\alpha} \right)$$

(2)

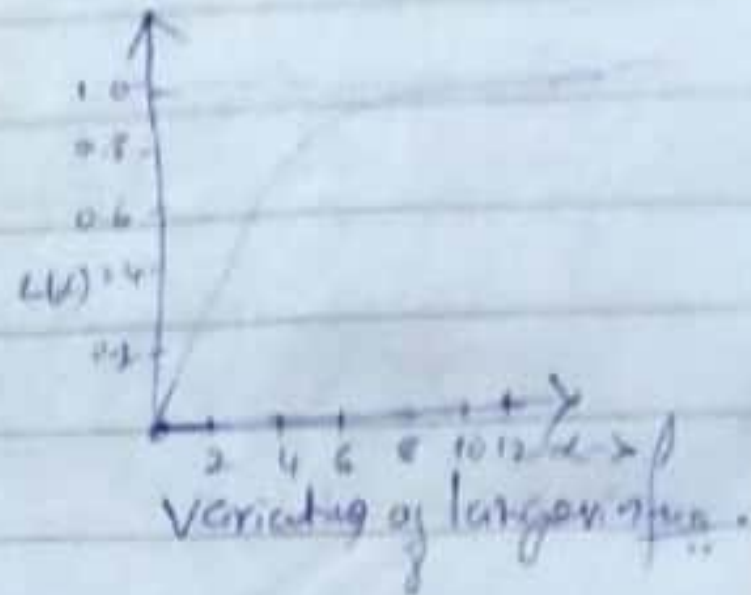
$$\textcircled{6} \quad \bar{I} = I_0 L(\alpha) \quad \rightarrow \textcircled{7}$$

where $L(\alpha) = \left[\coth \alpha - \frac{1}{\alpha} \right]$ is known as Langevin function, variation of Langevin function $L(\alpha)$ with α is as shown in fig. It is observed from the plot that, for large values of α the Langevin function tends to unity.

$\bar{I} = I_0$. In this case all atomic dipoles are parallel to \vec{B} and magnetisation attains saturation. For a small value of α the curve is linear and for smaller values of α the Langevin function becomes

$$L(\alpha) = \coth \alpha - \frac{1}{\alpha} = \alpha/3$$

Using α values $L(\alpha) = \frac{\mu B}{3kT}$



Using this value in Eqⁿ $\textcircled{7}$ $\bar{I} = I_0 \mu B / 3kT$ but $I_0 = Nm$

$$\therefore \bar{I} = \frac{Nm^2 B}{3kT} \quad \rightarrow \textcircled{8}$$

For paramagnetic substance, \bar{I} is very small positive value, hence $B = \mu_0(H + \bar{I}) = \mu_0 H$

Using this value in Equation $\textcircled{8}$

$$\therefore \bar{I} = \frac{Nm^2 \mu_0 H}{3kT} \quad \rightarrow \textcircled{9}$$

By the definition of magnetic susceptibility

$$\chi = \frac{\bar{I}}{H} = \frac{Nm^2 \mu_0 H}{\mu_0 H 3kT} = \frac{\mu_0 N m^2}{3kT}$$

$$\chi = \frac{\mu_0 M_0^2}{3kT} = \text{constant} = \frac{C}{T}$$

$$\chi = C/T \quad \rightarrow \textcircled{10}$$

hence $C = \mu_0 M_0^2 / 3k$ is a constant called Curie constant. Equation $\textcircled{10}$ is known as Curie law. It states that, the magnetic susceptibility of a paramagnetic substance varies inversely with the absolute temperature.

⑧ Hysteresis:

The term hysteresis is derived from a Greek word meaning "lagging behind". The lagging of intensity of magnetisation behind the magnetic field is known as hysteresis.

⑨ Hysteresis Curve:

Consider an unmagnetized ferromagnetic material. If it is subjected to a gradually increasing magnetizing field H . When $H=0$, $I=0$ i.e. origin O . As H increases gradually, I also increases in a non-linear way and attains a saturation at A as shown in fig. This is represented by a curve OA called the initial magnetization curve. At this stage, all the dipoles are aligned and I has reached to its maximum (or) saturated value.

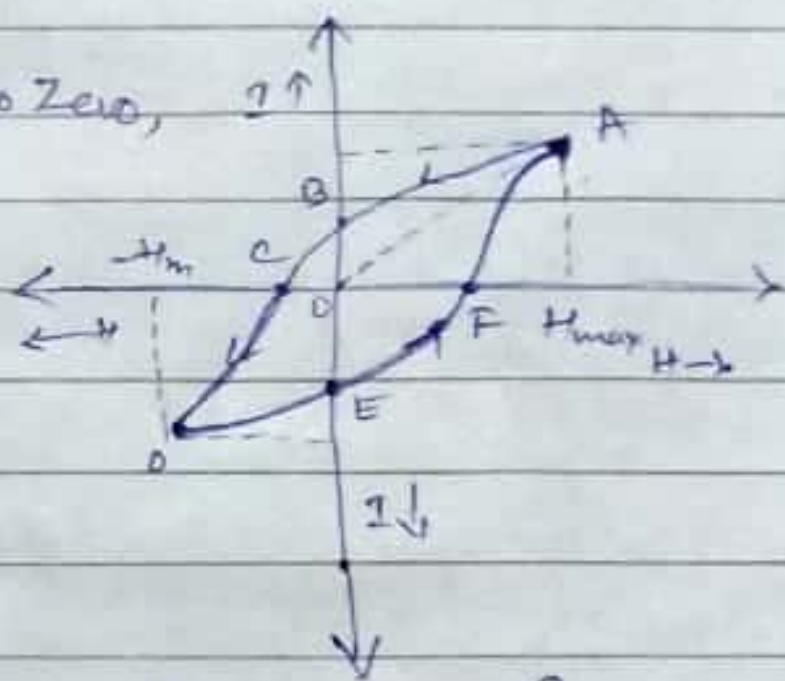
As H decreases from H_{max} to zero, it is found that the value of I also decreases but at a lower rate and follows the path AB .

When H is reduced to zero,

I is not come to its zero value.

but it has certain value in it

and is represented by OB .



Hysteresis Curve

The magnetization left behind in the specimen even when the magnetizing field is made zero is called residual magnetization

⑩ remanence. It gives the state of permanent magnetization of the specimen.

On applying the reverse field, the value of I falls and finally becomes zero. The abscissa OC represents the reverse magnetizing field needed to reduce magnetization to zero (or) to completely demagnetize the specimen and the field is known as coercive field. If the reverse magnetizing field is further increased, the magnetization increases in the reverse

④

direction and attain saturation at D. Again if the field is reduced to zero, I follows the path DE & remains equal to OE when $H=0$, here OE represents the residual magnetism.

When the field is again increased from zero to H_{max} , the curve EFA is traced and the saturation is again obtained in the +ve direction at A. The closed loop ABCDEFA represents the hysteresis curve of the material. It is observed that throughout the cycle of magnetization, I lags behind H. This lagging behind of I with respect to H is known as hysteresis.

② Significance of Hysteresis Curve :-

- It provides information regarding to the magnetic properties of the substance and helps to select suitable magnetic materials for various practical purposes.
- Hysteresis loops are different for different materials and one can compare the properties of the materials from the hysteresis loop.
- The area of the hysteresis loop gives the amount of energy lost during one cycle of magnetization.
- The choice of magnetic materials for different uses can be decided from the hysteresis curve of a specimen of the material.

③ Problems :-

- ① A bar magnet made of steel has a magnetic moment of 2.5 Am^2 and a mass of $6.6 \times 10^3 \text{ kg}$. If the density of steel is $7.9 \times 10^3 \text{ kg m}^{-3}$, find the intensity of magnetization of the magnet.

Given data: Magnetic moment (M) = 2.5 Am^2

mass (m) = $6.6 \times 10^3 \text{ kg}$

density (ρ) = $7.9 \times 10^3 \text{ kg m}^{-3}$

Horizon

⑤

Soln:- We have, the intensity of magnetisation $I = \frac{M}{V}$
where V is the volume of material.

$$V = \frac{\text{mass}}{\text{density}}$$

$$V = \frac{6.6 \times 10^{-3}}{7.9 \times 10^3}$$

$$V = 8.32 \times 10^{-7} \text{ m}^3$$

$$I = \frac{M}{V}$$

$$I = 3 \times 10^6 \text{ Am}^{-1}$$

③ Concept Questions:-

① Is it possible to obtain an isolated north or south pole by breaking a magnet into two pieces? why?

Ans: No.

Because according to Gauss law of magnetism isolated magnetic pole does not exist and magnetic flux through any closed surface is zero.

② Which of the two quantities related to magnetism have same dimensions?

Ans: Intensity of magnetisation and magnetising field intensity.

③ Can we have a magnetic hysteresis in paramagnetic or diamagnetic substance? why?

Ans: No.

because the hysteresis is a special property exhibited by only ferromagnetic materials.

④